

Multi-Electron Processes in MeV/u Mixed-State $C^{4+}(1s^2, 1s2s^3S) + He$ Collisions: Comparison of AOCC Calculations and Experiments

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Abstract

Complex three-electron Atomic Orbital Coupled Channel (AOCC) calculations using the semi-classical close-coupling approach [1] have been performed and compared to state-selective zero-degree Auger projectile spectroscopy measurements [2] for 4-18 MeV $C^{4+}(1s^2, 1s2s^3S)$ mixed-state ions in collisions with He. Single differential cross sections for the production of the three-electron $C^{3+} 1s2s^2 2S$, $1s2s2p 4^2P$, $1s2p^2 2D$ states from either the $1s^2$ ground or $1s2s^3S$ metastable beam component are determined using a new experimental technique [3] that can be applied in a two-measurement approach using different amounts of $C^{4+}(1s2s^3S)$ metastable fraction.

One of the existing more severe disagreements is the ratio of cross sections $R_m = \sigma_m(1s2s2p 4^2P) / \sigma_m(1s2s2p 2^2P)$ produced by direct transfer to the $1s2s^3S$ metastable (m) component [4-6], expected to be 2 according to spin statistics [7]. However, previous measurements [4,5], as well as theoretical calculations including cascade feeding (expected to have a significant role in selectively enhancing the $1s2s2p 4^2P$ population [4-6]) have found much larger numbers, in disagreement with the most recent results of $R_m \approx 2$ [8].

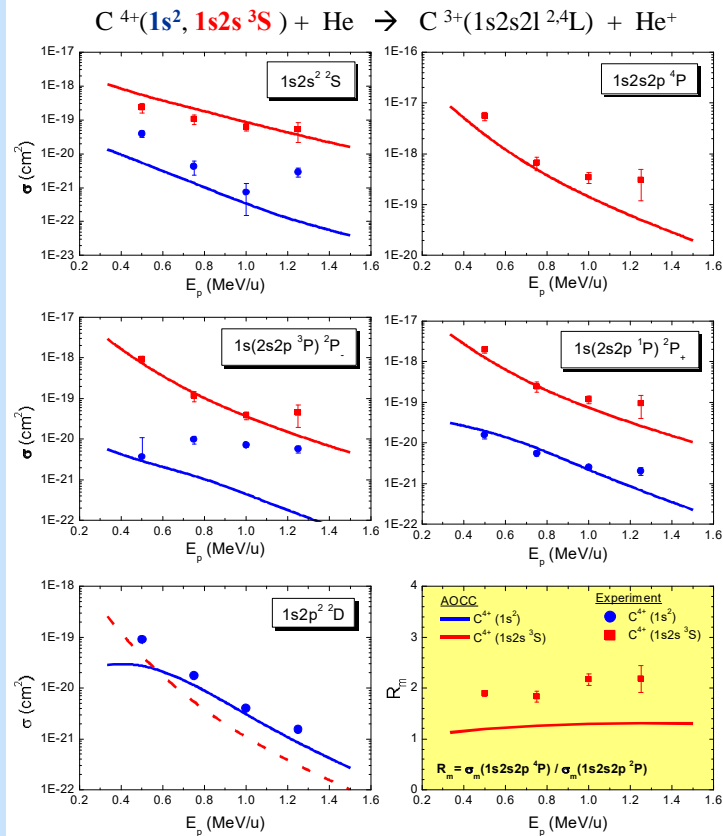


Fig. 2. Comparison of AOCC 3-electron calculations to experimental data obtained with zero-degree Auger projectile spectroscopy. AOCC calculations involve two projectile electrons and one target electron. The experimental SDCS data of Fig. 1 were multiplied with 4π to obtain the total cross sections. The graph in yellow background is the ratio $R_m = \sigma_m(4^2P) / \sigma_m(2^2P)$ that is examined in order to identify cascading contribution effects to the 4^2P .

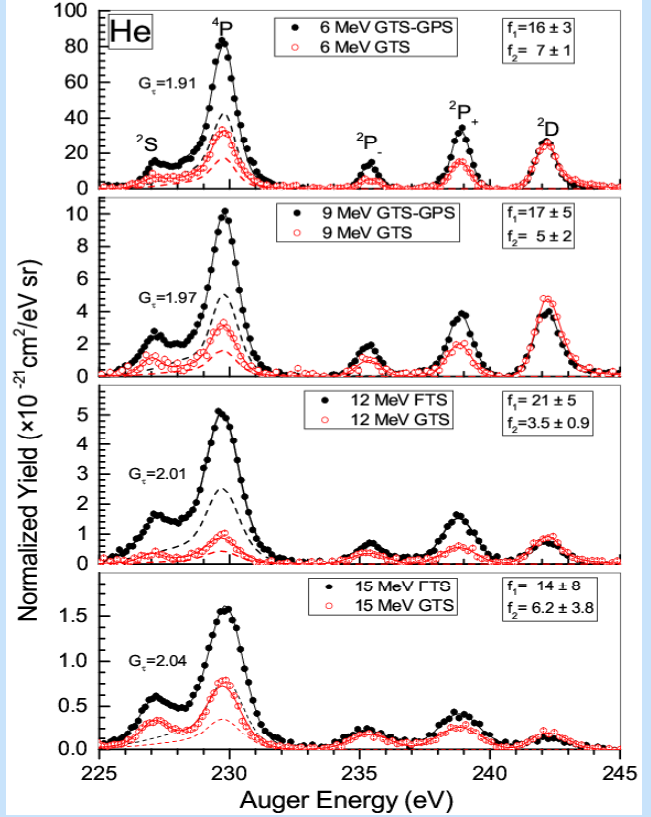


Fig. 1. Zero-degree C^{3+} Auger line spectra for $C^{4+}(1s^2, 1s2s^3S)$ collisions with He. Two spectra are shown for each collision energy with different metastable ($1s2s^3S$) beam fraction ($f_1 > f_2$) content, (in %) as judged by the corresponding 4^2P yields. These spectra are used within the method described in [3] to separate the contributions of the ground $C^{4+}(1s^2)$ and metastable $C^{4+}(1s2s^3S)$ parts to the formation of the $1s2n1^2 2^4L$ states.

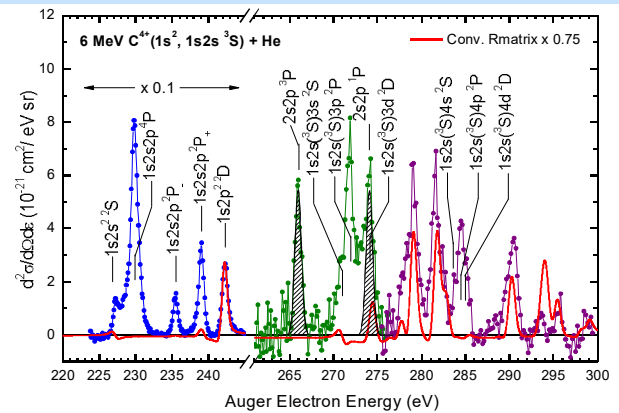


Fig. 3. Measured spectrum covering the whole KLn ($n \geq 2$) range. R-matrix calculations [9] convoluted with the spectrometer function account for the contributions from the $C^{4+}(1s^2)$. Capture to the $C^{4+}(1s2s^3S)$ lead to the populations of $1s2sn1^2 2^4L$, of which only the 2^4L strongly autoionize and can be seen in the spectrum, while the 4^4L cascade down to the $1s2s2p 4^2P$ state.

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